

$$\boxed{\vec{H} = -\nabla V_m}$$

This is possible only when  $\vec{J} = 0$ .

$\Rightarrow$  The scalar potential satisfies Laplace's equation :  $\boxed{\nabla^2 V_m = 0}$   
 $(\vec{J} = 0)$

Since,  $\nabla \cdot \vec{B} = 0 \rightarrow$  Non-existence of magnetic monopole.

$$\mu_0 \nabla \cdot \vec{H} = 0 \quad \therefore \vec{B} = \mu_0 \vec{H} \dots \text{free space}$$

$$\mu_0 \nabla \cdot (-\nabla V_m) = 0 \quad \therefore \vec{H} = -\nabla V_m$$

$$\boxed{\nabla^2 V_m = 0} \quad \text{only for } \vec{J} = 0$$

Directly defined as  $V_m = - \int_A^B \vec{H} \cdot d\vec{l} \dots (\text{unit: Ampere})$

$$\vec{H} \approx A/m$$

$$d\vec{l} \approx m.$$

$\Rightarrow$  Vector magnetic potential  
 $\hookrightarrow$  Exist where  $\vec{J}$  is present

" Magnetic vector potential is defined in such a way that its curl gives the magnetic flux density.

Using curl because  
curl of vector  $\approx$  Vector quantity

$$\boxed{\vec{B} = \nabla \times \vec{A}}$$

where  $\vec{A} \approx$  magnetic Vector potential  
 $(Wb/m)$

Since,  $\nabla \times \vec{H} = \vec{J} \dots$  Amper's circuit law.

$$\therefore \vec{B} = \mu_0 \vec{H} \dots \text{Free space}$$

$$\therefore \vec{H} = \frac{\vec{B}}{\mu_0}$$

$$\therefore \nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\text{Since, } \vec{B} = \nabla \times \vec{A}$$

$$\nabla \times \vec{B} = \nabla \times \nabla \times \vec{A} = 4\mu_0 \vec{J}$$

$$\therefore \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A} = 4\mu_0 \vec{J}$$

↓  
= 0

Since, For D.C. current only

$$\nabla \cdot \vec{A} = 0$$

(because magnetic field is continuous and it has no divergence.)

using Laplacian of Vector

$$\text{i.e. } \nabla^2 \vec{A} = \nabla \cdot (\nabla \cdot \vec{A}) - \nabla \times \nabla \times \vec{A}$$

$$\nabla \times \nabla \times \vec{A} = \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$$

$$-\nabla^2 \vec{A} = 4\mu_0 \vec{J}$$

$$\text{or } \nabla^2 \vec{A} = -4\mu_0 \vec{J}$$

↳ Vector Poisson's equation.

$$\nabla^2 \vec{A} = -4\mu_0 \vec{J} \quad \text{for } (x, y, z) \text{- Cartesian coordinate system.}$$

$$\nabla^2 (A_x \vec{a}_x + A_y \vec{a}_y + A_z \vec{a}_z) = -4\mu_0 (J_x \vec{a}_x + J_y \vec{a}_y + J_z \vec{a}_z)$$

$$\left. \begin{array}{l} \nabla^2 A_x = -4\mu_0 J_x \\ \nabla^2 A_y = -4\mu_0 J_y \\ \nabla^2 A_z = -4\mu_0 J_z \end{array} \right\}$$

These are the Poisson's equation in magnetostatics.

⇒ In case of Poisson's equation in Electrostatics

$$\nabla^2 V = -\frac{\rho_v}{\epsilon_0}$$

$$V = \frac{Q}{4\pi\epsilon_0 r}$$

volumic charge density  
↓  
 $\therefore Q = \int \rho_v dv$

$$A_x = \mu_0 J_x \left[ \frac{1}{4\pi} \int_V \frac{dv}{r} \right]$$

$$A_x = \frac{\mu_0}{4\pi} \int_V \left( \frac{J_x}{r} \right) dv$$

$$A_y = \frac{\mu_0}{4\pi} \int_V \left( \frac{J_y}{r} \right) dv$$

$$A_z = \frac{\mu_0}{4\pi} \int_V \left( \frac{J_z}{r} \right) dv$$

$$V = \frac{1}{4\pi\epsilon_0} \int_V \frac{\rho_v dv}{r}$$

$$= \frac{\rho_v}{\epsilon_0} \left[ \frac{1}{4\pi\epsilon_0} \int_V \frac{dv}{r} \right]$$

In general

$$\boxed{\vec{A} = \frac{\mu_0}{4\pi} \int_V \frac{\vec{J}}{r} dv}$$

$$\text{with } \vec{B} = \nabla \times \vec{A}$$

$$\nabla \cdot \vec{B} = 0$$

$$\therefore \nabla \cdot (\nabla \times \vec{A}) = 0$$

$\Rightarrow$   $\vec{A}$  in terms of three standard current configuration

i.e.  $\boxed{\int d\vec{l} = \vec{k} ds = \vec{J} dv} \rightarrow \text{Unit (Ampere-meter)}$

so,  $\vec{A} = \frac{\mu_0}{4\pi} \int_V \frac{\vec{J} dv}{r} \dots \text{Volume current}$

$$\vec{A} = \frac{\mu_0}{4\pi} \int_S \frac{\vec{k} ds}{r} \dots \text{Surface current}$$

$I = \text{Ampere}$

$d\vec{l} = \text{meter}$

$$\vec{k} = \frac{I}{b} \text{ (A/m)}$$

$ds = \text{meter}^2$

$$\vec{J} = \frac{I}{s} \text{ (A/m}^2\text{)}$$

$$dv = \text{m}^3$$